

Detonation in TATB Hemispheres

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Propellants, Explosives, Pyrotechnics

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Detonation in TATB Hemispheres

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Abstract

Streak camera breakout and Fabry-Perot interferometer data have been taken on the outer surface of 1.80 g/cm 3 TATB hemispherical boosters initiated by slapper detonators at three temperatures. The slapper causes breakout to occur at 54 $^\circ$ at ambient temperatures and 42 $^\circ$ at –54 $^\circ$ C, where the axis of rotation is 0 $^\circ$. The Fabry velocities may be associated with pressures, and these decrease for large timing delays in breakout seen at the colder temperatures. At room temperature, the Fabry pressures appear constant at all angles. Both fresh and decade-old explosive are tested and no difference is seen. The problem has been modeled with reactive flow. Adjustment of the JWL for temperature makes little difference, but cooling to -54° C decreases the rate constant by $1/6^{th}$. The problem was run both at constant density and with density differences using two different codes. The ambient code results show that a density difference is probably there but it cannot be quantified.

Keywords: booster, slapper detonator, Fabry, streak camera, reactive flow

1 Experimental

Hemispherical TATB boosters are often used to ignite larger charges of LX-17. In our previous paper,[1] we lit the ambient 1.80 g/cc ultrafine TATB with a slapper detonator, as shown in the schematic of **Figure 1**. The X-axis is the axis of rotation so that there is cylindrical symmetry. The LX-16 detonator drives an aluminum slapper across the air space to hit the TATB in part of section 1. The detonation spreads into sections 2 and 3 and finally to the quarter-circular edge where measurements are made. The angles from the origin are shown, with 0° lying on the axis of revolution. As a result of the asymmetric slapper initation, first breakout of the detonation at the outer surface of the 19 mm-radius booster occurred at roughly 60°. Computer modeling of the problem was difficult. The initial radial zoning was abandoned because of the rectangular symmetry of the slapper and reduced connectivity meshing was used, as shown by the sections in Figure 1. Program burn gave a breakout at 80° and definitely did not work. Reactive flow gave good results but was sensitive to all the input parameters. Possible density variations in the TATB made the problem even more difficult. Better data and analysis is offered here.

The geometry of Figure 1 was again used. The slapper detonator and its effect on TATB have previously been described. [2] The measured times are referenced to the bridge-wire burst that initiates the LX-16. Five curved disks of aluminum foil and Lucite of 8 mm diameter were placed so that a five-beam Fabry-Perot interferometer (Fabry) could measure the foil velocity at 3, 30, 60, 75 and 85 degrees.

Hemispheres were measured at ambient, -20°C and -54°C. Two kinds of TATB were used: "fresh" and 8-13 years older.

Figure 2 shows measured streak camera data in terms of bridgewire-burst times. As the sample cools, the curves tighten up with the first breakout angles becoming smaller. We have analyzed the data more thoroughly and show the results in **Table 1**. The averages at the bottom cover old and new as there is no difference within error. The breakout angle decreases from $54^{\circ} \pm 4$ at ambient temperatures to $42^{\circ} \pm 4$ at -54° C, which is a definite change of position. The average breakout angle is the same within error at both -20° and -54° C. The time from the first breakout to time of breakout at 0° on the axis is also given. because the error is only ± 0.01 μs. This time decreases from 0.09 μs at ambient to 0.06 μs at both cold temperatures. The bridgewire-burst times to first breakout are the same at 3.87, 3.82 and 3.82 μs, all ± 0.05 μs from ambient to coldest.

For the purposes of modeling, we need the averages of all these runs, and these are listed in **Table 2**. The three temperatures are shown along with the high and low bounds. These bounds will be used to illustrate the code runs.

The Fabry shows a rapid jump-up to a maximum velocity value, which then decays slowly with a reshock often pushing it up again. The reshocks appear to be caused by reflections from the sides of the small square foil/LiF packages. The maximum value of particle velocity is roughly related to the maximum pressure in the plastic by the equation

$$P = \rho_o \left(C_o + S_I u_p \right) u_p \tag{1}$$

where the density of lucite is 1.182 g/cc, C_0 is 2.18 mm/ μ s and $S_1 = 2.088$.

The Fabry maximum values of measured velocity and calculated pressure are listed in **Table 3**. The scatter in the Fabry data can be estimated from four sets of repeated points, where sets are defined by temperature and age. With this limited calibration, we estimate the measured velocity error to be ± 0.1 mm/ μ s or ± 1.6 GPa.

It is of interest to try to correlate the shape of the breakout time curves with the pressures obtained from the measured Fabry velocities. These are shown in **Figure 3** at ambient temperatures, where there are only two old runs and and one new one. The new data lies lower, but, using the error bars derived above, the curves must be judged the same. It is not possible to unravel pressure-angle information from ambient data at this time.

We now expand out to all the data in **Figure 4**. In each set, we calculate the pressure difference relative to the 60° point for both the pressure and the breakout time. There is indeed a linear relationship caused by the large angle data at cold temperatures. The width of the line is about ± 3 GPa so that the finer relations at small angles cannot be resolved.

Figure 5 shows the results of plotting the ratio of new to old material velocities at each angle. The ambient data is a little low, but this is the result of the single new set. Only at 85° is there some evidence of possible change. Overall, a decade of age appears not to have affected the TATB properties.

2 Calculations

We reran this problem in JWL++ using the reduced connectivity mesh shown as regions in Figure 1. Section 1 is a square grid at 7 zones/mm used to match the slapper. Sections 2 and 3 are semi-radial in nature. Although not shown here accurately, we added the major air gaps around the LX-16 driver and the foil on the front of it so that the initiation system has the most accurate depiction to date. We use the simple reactive flow model, JWL++, inside a 2-dimensional arbitrary Langragian-Eulerian (ALE) hydrocode of the CALE-type.[3] The detonation rate is given by a single rate constant, G₁, using

$$\frac{dF}{dt} = G_I(P+Q)^{b_I} (1-F) \approx \frac{1}{2} G_I P_{cj}^{b_I}$$
(2)

where the second term is what the code gets and the third is an analytical approximation. Decreasing the rate constant slows down the detonation velocity and lengthens the reaction zone. This model creates a dead zone along the edges of the explosive, but detonation failure is not included. The model fails occurs at $b_1 = 3$ so that increasing b_1 increases the potential for failure.[4] However, increasing b_1 makes the size (diameter) effect curve more concave down, and the data is slightly concave-up. So a b_1 of either 1 or 2 is usually used. Running near $b_1 = 3$ makes the model unstable.

An important part of our experiment are the cooled samples, and we want to create a JWL equation of state that corrects for contraction and energy change upon becoming colder. For changing the temperature without having a real temperature in a JWL, we use as input the coefficient of volumetric expansion,

 \tilde{V} in K^{-1} , and the heat capacity at constant pressure per unit volume, C_p , both functions of the absolute temperature, T. We have

$$\Delta V / V_{O}(dim \ ensionless) = 10^{-6} (\beta_{I} + \beta_{2}T + \beta_{3}T^{2}) C_{p}(kJ/cm^{3} \cdot K) = 10^{-3} (C_{pI} + C_{p2}T)$$
(3)

Both are integrated from the reference temperature of 25°C (298.15 K) to give the change of volume and energy between room temperature and T:

$$\Delta v(dim\ ensionless\) = 10^{-6} \left(\beta_I (T - 298.15\) + \frac{\beta_2}{2} (T^2 - 298.15^2\) + \frac{\beta_3}{3} (T^3 - 298.15^3\) \right)$$

$$\Delta E(kJ/cm^3\) = 10^{-3} \left(C_{pI} (T - 298.15) + \frac{C_{p2}}{2} (T^2 - 298.15^2\) \right)$$
(4)

The original room temperature density is converted to the new density by

$$\rho_O(new) = \frac{\rho_O}{I + \Delta v}.$$
 (5)

The detonation energies are changed by

$$E_d(new) = \frac{E_d(old) + \Delta E}{I + \Delta v} \tag{6}$$

As a means of estimating how the detonation energies might change, we use this relationship

$$U_{S}(new) \approx \left(\frac{E_{O}(new)}{E_{O}(old)}\right)^{1/2} U_{S}(old) \tag{7}$$

where we assume that the full E_0 is dumped into the explosive immediately even though it is not available for external work until the products expand to infinite volume. For TATB, the values for β_1 through β_3 are 551, -3.84 and 0.00993 in dimensionless units.[5] The values of C_{p1} and C_{p2} in 0.529 and 0.0044 J/cm³·K.[5,6] The resulting JWL has a higher energy density because of contraction but energy is lost from detonation to make up in the cooling. In cooling TATB to -54°C, the density increases from 1.80 to 1.833 g/cc, but the detonation energy at v = 2.2 changes from 4.44 to 4.45 kJ/cc and the detonation velocity from 7.59 to 7.592 mm/ μ s.

In **Figure 6**, we plot the measured data ranges as listed in Table 2. The broad dark gray lines indicate the range at ambient temperatures and the light gray -54° C. The curve passing through the ambient range is $b_1 = 2$ with $G_1 = 0.06$ (μ s. GPa^2)⁻¹. There are two $b_1 = 2$ curves passing through the low temperature data. The left one uses the same room temperature JWL as the ambient run. The right one has the temperature-adjusted JWL. There is little change obtained by adjusting the JWL, but both have G_1 lowered to 0.05 (μ s· GPa^2)⁻¹ to fit at -54° C. We thus have evidence that the rate constant decreases with temperature. If we run at room temperature with $b_1 = 1$, which is a reasonable setting, the curve lies too far to the right with a break-out angle of 60- 70° . This suggests that we need failure mixed it somehow in this problem and we are getting it here by raising b_1 . We note, however, that the fits are not so good at small angles.

As reported before, the boosters were ram-pressed and X-ray tomography showed a variation of density.[1] The absolute calibration of this effect has not yet been possible, and we overestimated the differences last time. We may roughly align Figure 1 in the model so that section 1 remains at 1.80 g/cm^3 , section 2 is higher at 1.825 g/cm^3 and section 3 is lower, at 1.775 g/cm^3 . The ambient result for $b_1 = 2$, $G_1 = 0.06 (\mu \text{s.GPa}^2)^{-1}$ is shown in **Figure 7**. Fitting the large-angle part is easy but the small-angle part changes. This happens because the on-axis section has a lower detonation velocity than expected, so that the breakout along the axis occurs later. The all- 1.80 g/cm^3 result is too low in this region, as noted before, but the $1.775/1.825 \text{ g/cm}^3$ run is too high. It is possible that a split of 1/3 to 1/2 is closer to the right answer.

It seems important to move from two to three dimensions in the modeling, even though resolution is hard to get when adding the third dimension. We also ran the problem using Ignition & Growth in 3-dimensional LSDYNA, which is also an ALE code.[7, 8] Again, we had a 5 mm center square section at 8 zones/mm, which changed into a graduated mesh beyond. A total of 510870 nodes and 486692 elements were included. A parabolic shaped aluminum flyer moving at 3.5 mm/µs simulated the detonator. A 1.5 mm thick lucite shell was placed around the outer edge of the TATB to simulate the measuring apparatus. Ignition & Growth is a more complicated reactive flow code requiring twice the minimum zoning as JWL++, so that the 8 zones/mm are needed. Its rate equation contains the form $F^{0.667}(1-F)P^2$ with a fast-reaction rate constant of $0.22 \, (\mu s \, GPa^2)^{-1}$.[9] The results are shown in Figure 8. Again, the small-angle region is raised by the decrease in density in this section. This run suggests that a spread as large as $1.825/1.775 \, g/cm^3$ may be correct.

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Symbols and Abbreviations

- b₁ Pressure coefficient in the rate (dimensionless)
- C_o U_s - u_p coefficient (mm/ μ s)
- C_p Heat capacity (kJ/cm³·K)
- C_{p1} Heat capacity coefficient (kJ/cm³ K)
- C_{p2} Heat capacity coefficient (kJ/cm³)
- E_d Detonation energy at some relative volume (kJ/cm³)
- E_0 Total detonation energy (kJ/cm³)
- ΔE Change in detonation energy (kJ/cm³)
- F Burn fraction (dimensionless)
- G_1 Detonation rate constant $(\mu s GPa^{b1})^{-1}$
- P Pressure (GPa
- P_{cj} C-J pressure (GPa)
- Q Artificial viscosity (GPa)
- S_1 U_s - u_p coefficient (dimensionless)
- T Temperature (K)
- t Time (µs)
- U_s Explosive detonation velocity (mm/μs)
- u_{p} Explosive particle velocity (mm/ μ s)
- V₀ Initial Volume (cm³)
- ΔV Change in volume (cm³)
- Δv Change in relative volume (dimensionless)
- β₁ Volume expansion coefficient (dimensionless)

- Volume expansion coefficient (K⁻¹) Volume expansion coefficient (K⁻²) β_2
- β_3

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Table 1. Streak camera summary of all shots with the TATB hemispheres. Old and new refers to the age of the TATB.

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new 45 3 0.06 3.780 -20 8/28/01 39 2 0.05 3.730 old 43 3 0.06 3.720 -20 9/4/02 46 3 0.07 3.770 old 42 3 0.06 3.780 -54 4/12/01 45 5 0.05 3.730 -54 9/10/02 40 3 0.07 3.790 new 42 3 0.06 3.800 -54 8/17/01 43 4 0.06 3.720 old 40 5 0.05 3.730 -54 8/31/01 42 4 0.06 3.730 -54 8/31/01 42 4 0.06 3.730 -54 9/17/02 42 4 0.05 3.860 -54 9/17/02 42 4 0.05 3.860 old 42 5 0.0		old	63	7	0.08	3.690
-20 8/28/01 39 2 0.05 3.730 old 43 3 0.06 3.720 -20 9/4/02 46 3 0.07 3.770 old 42 3 0.06 3.780 -54 4/12/01 45 5 0.05 3.730 new 37 4 0.05 3.730 new 42 3 0.06 3.800 -54 8/17/01 43 4 0.06 3.720 old 40 5 0.05 3.730 old 40 5 0.05 3.730 -54 8/31/01 42 4 0.06 3.730 old 43 5 0.05 3.740 old 43 5 0.05 3.740 old 42 5 0.06 3.850	-20	9/16/02	50	4	0.06	3.780
old 43 3 0.06 3.720 -20 9/4/02 46 3 0.07 3.770 old 42 3 0.06 3.780 -54 4/12/01 45 5 0.05 3.730 -54 9/10/02 40 3 0.07 3.790 new 42 3 0.06 3.800 -54 8/17/01 43 4 0.06 3.720 old 40 5 0.05 3.730 -54 8/31/01 42 4 0.06 3.730 -54 8/31/01 42 4 0.06 3.730 -54 9/17/02 42 4 0.05 3.860 -54 9/17/02 42 4 0.05 3.860 old 42 5 0.06 3.850 ambient average 54 4 0.09 3.785		new	45	3	0.06	3.780
-20 9/4/02 old 46 old 3 old 0.07 old 3.770 old -54 4/12/01 d5 old 5 old 3.730 old 3.730 old -54 9/10/02 d0 old 3 old 3.730 old 3.730 old -54 9/10/02 d0 old 3 old 3.790 old 3.800 old -54 8/17/01 d3 old 4 old 5 old 3.720 old old d40 old 5 old 3.730 old 3.730 old -54 8/31/01 old 42 old 4 old 3.730 old -54 9/17/02 d2 d2 old 4 old 3.860 old -54 9/17/02 d2 d2 old 4 old 3.850 old ambient average 54 description 4 old 3.785	-20	8/28/01	39	2	0.05	3.730
old 42 3 0.06 3.780 -54 4/12/01 45 5 0.05 3.730 new 37 4 0.05 3.730 -54 9/10/02 40 3 0.07 3.790 new 42 3 0.06 3.800 -54 8/17/01 43 4 0.06 3.720 old 40 5 0.05 3.730 -54 8/31/01 42 4 0.06 3.730 old 43 5 0.05 3.740 -54 9/17/02 42 4 0.05 3.860 old 42 5 0.06 3.850 ambient average 54 4 0.09 3.785		old	43	3	0.06	3.720
-54 4/12/01 45 5 0.05 3.730 new 37 4 0.05 3.730 -54 9/10/02 40 3 0.07 3.790 new 42 3 0.06 3.800 -54 8/17/01 43 4 0.06 3.720 old 40 5 0.05 3.730 -54 8/31/01 42 4 0.06 3.730 old 43 5 0.05 3.740 -54 9/17/02 42 4 0.05 3.860 old 42 5 0.06 3.850 ambient average 54 4 0.09 3.785	-20	9/4/02	46	3	0.07	3.770
new 37 4 0.05 3.730 -54 9/10/02 40 3 0.07 3.790 new 42 3 0.06 3.800 -54 8/17/01 43 4 0.06 3.720 old 40 5 0.05 3.730 -54 8/31/01 42 4 0.06 3.730 old 43 5 0.05 3.740 old 43 5 0.05 3.740 -54 9/17/02 42 4 0.05 3.860 old 42 5 0.06 3.850 ambient average 54 4 0.09 3.785		old	42	3	0.06	3.780
-54 9/10/02 40 3 0.07 3.790 new 42 3 0.06 3.800 -54 8/17/01 43 4 0.06 3.720 old 40 5 0.05 3.730 -54 8/31/01 42 4 0.06 3.730 old 43 5 0.05 3.740 -54 9/17/02 42 4 0.05 3.860 old 42 5 0.06 3.850 ambient average 54 4 0.09 3.785	-54	4/12/01	45	5	0.05	3.730
new 42 3 0.06 3.800 -54 8/17/01 43 4 0.06 3.720 old 40 5 0.05 3.730 -54 8/31/01 42 4 0.06 3.730 old 43 5 0.05 3.740 -54 9/17/02 42 4 0.05 3.860 old 42 5 0.06 3.850 ambient average 54 4 0.09 3.785		new	37	4	0.05	3.730
-54 8/17/01 43 4 0.06 3.720 old 40 5 0.05 3.730 -54 8/31/01 42 4 0.06 3.730 old 43 5 0.05 3.740 -54 9/17/02 42 4 0.05 3.860 old 42 5 0.06 3.850 ambient average 54 4 0.09 3.785	-54	9/10/02	40	3	0.07	3.790
old 40 5 0.05 3.730 -54 8/31/01 42 4 0.06 3.730 old 43 5 0.05 3.740 -54 9/17/02 42 4 0.05 3.860 old 42 5 0.06 3.850 ambient average 54 4 0.09 3.785		new	42	3	0.06	3.800
-54 8/31/01 42 4 0.06 3.730 old 43 5 0.05 3.740 -54 9/17/02 42 4 0.05 3.860 old 42 5 0.06 3.850 ambient average 54 4 0.09 3.785	-54	8/17/01	43	4	0.06	3.720
-54 8/31/01 42 4 0.06 3.730 old 43 5 0.05 3.740 -54 9/17/02 42 4 0.05 3.860 old 42 5 0.06 3.850 ambient average 54 4 0.09 3.785		old	40	5	0.05	3.730
old 43 5 0.05 3.740 -54 9/17/02 42 4 0.05 3.860 old 42 5 0.06 3.850 ambient average 54 4 0.09 3.785	-54	8/31/01	42	4	0.06	
-54 9/17/02 42 4 0.05 3.860 old 42 5 0.06 3.850 ambient average 54 4 0.09 3.785				5		
old 42 5 0.06 3.850 ambient average 54 4 0.09 3.785	-54					
ambient average 54 4 0.09 3.785						
_	ambient					
-	-20	_	44	3	0.06	3.760
	-54	· ·	42	4	0.06	3.768

Table 2. Breakout time limits in μs for the three temperatures. The high and low bounds are listed.

Angle	Ambient		-20°C		-55°C	
(degrees)	High	Low	High	Low	High	Low
0	0.10	0.08	0.07	0.06	0.08	0.05
15	0.09	0.07	0.06	0.04	0.05	0.03
30	0.05	0.03	0.03	0.02	0.02	0.00
45	0.02	0.00	0.01	0.00	0.01	0.00
60	0.01	0.00	0.08	0.01	0.11	0.03
75	0.09	0.05	0.23	0.12	0.39	0.20
90	0.28	0.16	0.54	0.41	0.80	0.48

Table 3. Summary of Fabry measurements. Old and new refers to the age of the TATB.

14010 01 2	anning of	r dory meds	Maximum	Maximum	Peak
Temp.	Shot	Angle	Velocity		Times
(°C)	Name	(degrees)	(mm/µs)	(GPa)	(μs)
ambient	09/11/02	3	2.71	25.1	3.91
amorem	new	30	2.77	26.1	3.87
	IIC W	60	2.69	24.8	3.78
		75	2.72	25.2	3.86
		85	2.72	25.0	3.92
ambient	1/27/00	7	2.95	29.1	3.92
amorem	old	30	2.90	28.2	3.89
	olu	60	2.77	26.0	3.84
		75	2.83	27.0	3.88
		85	2.65	24.2	3.98
ambient	1/28/00	7	2.84	27.2	3.90
amorem	old	30	2.88	27.9	3.86
	Olu	60	2.85	27.4	3.82
		75	2.91	28.4	3.86
		85	2.64	24.0	3.94
-20	9/16/02	3	2.77	26.1	3.92
20	new	30	2.65	24.1	3.87
	IIC W	60	2.71	25.1	3.85
		75	2.71	18.6	4.05
		85	2.41	20.5	4.14
-20	8/28/01	3	2.71	25.1	3.83
-20	old	30	2.70	24.9	3.79
-20	9/4/02	3	2.91	28.4	3.91
20	old	30	2.93	28.7	3.87
	014	60	2.60	23.4	3.89
		75	2.64	24.1	4.04
		85	2.38	20.1	4.18
-54	4/12/02	7	2.66	24.3	3.84
	new	30	2.57	22.9	3.81
		60	2.45	21.1	3.87
-54	9/10/02	3	2.77	26.1	3.90
	new	30	2.71	25.1	3.85
		60	2.73	25.4	3.94
		75	1.46	9.0	3.93
		85	1.83	13.0	4.43
-54	8/17/01	3	2.77	26.1	3.92
	old	30	2.65	24.1	3.87
		60	2.71	25.1	3.85
		75	2.27	18.6	4.05
		85	2.41	20.5	4.14
-54	8/31/101	3	2.71	25.1	3.87
	old	30	2.60	23.4	3.82
		60	2.60	23.4	3.91
-54	9/17/02	3	2.90	28.2	3.96
	old	30	2.74	25.5	3.90
		60	2.58	23.0	4.01

75	2.12	16.5	4.33
85	1.97	14.7	4.54

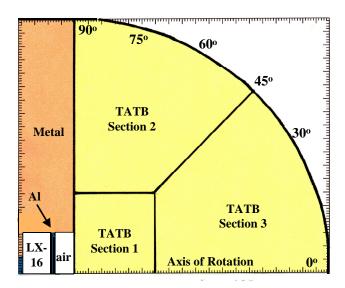


Figure 1. Schematic of the TATB booster, detonator and metal back-up plate. The X-axis is an axis of rotation. The detonation runs from the lower left up and to the right. Measurements are made on the quarter-circle edge at the angles listed in degrees.

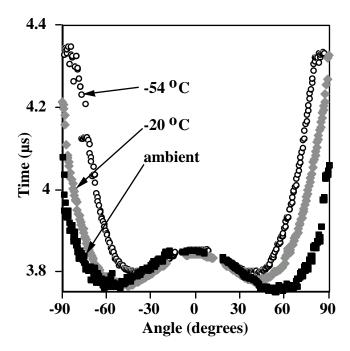


Figure 2. Measured streak camera data showing the first breakout in the vicinity of 60°. As the TATB becomes colder, the curve tightens up, with the first breakout moving to smaller angles.

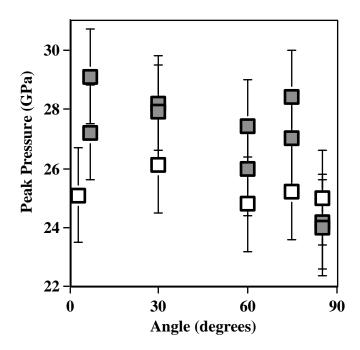


Figure 3. Calculated pressures from Fabry velocity data as a function of angle at ambient temperatures. New points are open and old points are gray.

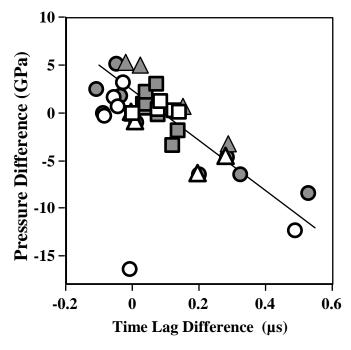


Figure 4. Pressure difference relative to the 60° pint that set versus the same breakout time difference. There is a linear relationship on a broad scale caused mostly by the large angles at cold temperatures. One point is way out of line.

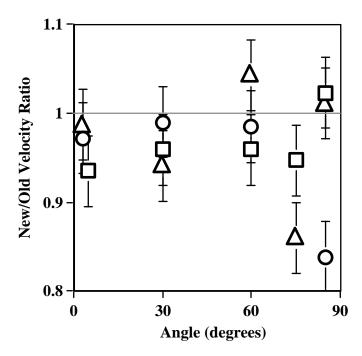


Figure 5. Ratio of measured velocities of new and old explosive as a function of angle. The temperatures are: ambient (square), -20° C (triangle) and -55° C (circle). It is difficult to see a trend in the data.

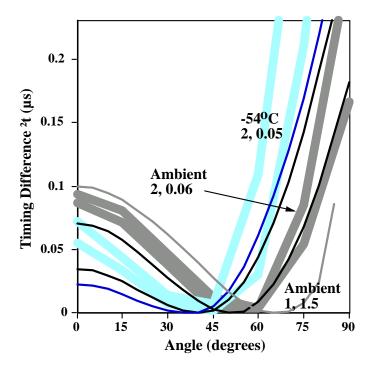


Figure 6. JWL++ runs with streak camera data ranges shown by the broad gray lines for room temperature and the light gray lines for -54° C. The code runs for $b_1 = 2$ are close to these ranges, whereas the $b_1 = 1$ run is far off. In order to make the model transition to cold temperature, the rate constant is changed from $0.06~(\mu s.GPa^2)^{-1}$ to 0.05 with cooling.

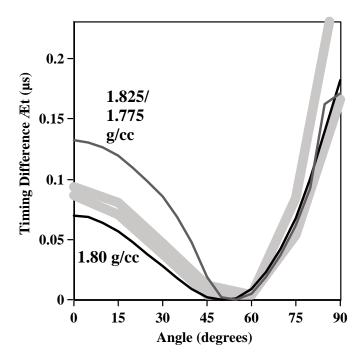


Figure 7. JWL++ ambient $b_1 = 2$ runs at constant and variable density. The actual density spread appears to lie in between.

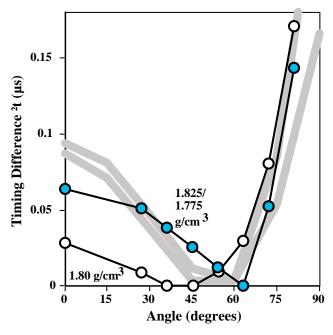


Figure 8. Ambient LSDYNA run in 3-dimensions showing the difference between constant and variable density.